

# Chlorine-NQR Relaxation Near the Structural Transitions of Natural and Deuterated $(\text{NH}_4)_2\text{TeCl}_6$ \*

C. Dimitropoulos<sup>1</sup>, F. Borsa<sup>2</sup>, and J. Pelzl<sup>3</sup>

<sup>1</sup> Inst. de Physique Expérimentale, EPFL, PHB-Ecublens, 1015 Lausanne, Switzerland

<sup>2</sup> Dipartimento di Fisica dell'Università e I.N.F.N., 27100 Pavia, Italy

<sup>3</sup> Inst. für Experimentalphysik VI, Ruhr-Universität, 4630 Bochum 1, FRG

Z. Naturforsch. **47a**, 261–264 (1992); received July 27, 1991

The temperature and isotope dependence of the  $^{35,37}\text{Cl}$ -NQR spin-lattice relaxation rate near the cubic to trigonal phase transition at  $T_{c1} \cong 87$  K has been investigated in both  $(\text{NH}_4)_2\text{TeCl}_6$  and  $(\text{ND}_4)_2\text{TeCl}_6$  salts. In the temperature range of  $\sim 20$  K above and  $\sim 10$  K below  $T_{c1}$ , the relaxation obeys the relation  $(T_1 \cdot T)^2 \propto (T - T_{c1})$ . This behaviour corresponds to a direct relaxation mechanism in presence of a strongly damped rotational soft mode of the  $\text{TeCl}_6^{2-}$  tetrahedra and is discussed on the basis of cation-anion coupling near the critical region. Below  $T_{c1}$  we observe an extra contribution to the  $^{35,37}\text{Cl}$  relaxation rate which is tentatively ascribed to the dynamics of the  $\text{NH}_4^+$  tetrahedra.

PACS numbers : 76.60.Gv, 76.60.Es, 64.60.Ht.

## Introduction

The natural and deuterated ammonium hexachlorotellurate  $(\text{NH}_4)_2\text{TeCl}_6$  has, at room temperature, the same cubic antiferroite crystal structure as the  $\text{A}_2\text{BX}_6$  compounds. This structure belongs to the space group  $\text{O}_h^5$  ( $\text{Fm}\bar{3}\text{m}$ ) and is characterized by a face-centered cubic arrangement of the  $\text{BX}_6^{2-}$  octahedra surrounded by a single cubic cage of  $\text{A}^+$  ions [1]. Normally, most of the  $\text{A}_2\text{BX}_6$  hexahalometallates undergo a displacive transition involving a tetragonal distortion accompanied, by the softening of a rotary phonon mode of the  $\text{BX}_6^{2-}$  octahedra [2]. On the other hand, in the  $(\text{NH}_4)_2\text{TeCl}_6$  and  $(\text{ND}_4)_2\text{TeCl}_6$  crystals studied here, a structural transition at  $T_{c1} \sim 85$  K, associated to a trigonal distortion was found from chlorine-NQR and Raman scattering [3–5]. The distinctive feature of this transition compared to the others occurring in  $\text{A}_2\text{BX}_6$  antiferroites, is that all halogens remain at equivalent sites below  $T_c$ . Thus, a single chlorine-NQR line is observed through the transition, providing an ideal system to investigate the critical dynamics above and below  $T_c$ . Furthermore, unusual isotope effects on its structural phase transition have been reported whereby, on deuteration two additional

transformations are induced below 80 K, which have been explained by a condensation of the deuteron motion in a triangular shaped potential [6, 7].

Here we report on a detailed study of the critical behavior of the  $^{35,37}\text{Cl}$ -NQR relaxation in the vicinity of  $T_{c1}$  aimed at investigating the effect of the coupling between the anion  $\text{TeCl}_6^{2-}$  octahedra and the cation  $\text{NH}_4^+$  tetrahedra on the rotary soft mode, with particular regard to the isotope dependence of this coupling.

## Experimental

Measurements of the NQR frequency and relaxation were performed using a modified Bruker SXP-100 spectrometer and an Oxford CF-200 continuous flow cryostat. The temperature of the samples was stabilized to better than 0.1 K. Single crystals and polycrystalline samples of  $(\text{NH}_4)_2\text{TeCl}_6$  and  $(\text{ND}_4)_2\text{TeCl}_6$  were prepared by precipitation in aqueous solutions. Deuterated crystals have been prepared using DCl and  $\text{D}_2\text{O}$ . The degree of deuteration has been monitored by Raman scattering from the internal ammonium ion modes.

## Results

The temperature dependence of the  $^{35}\text{Cl}$ -NQR frequency  $\nu_Q$  is shown in Fig.1 for the natural and deuterated ( $\sim 95\%$ ) compounds. In both cases, the

\* Presented at the XIth International Symposium on Nuclear Quadrupole Resonance Spectroscopy, London, U.K., July 15–19, 1991.

Reprint requests to Dr. C. Dimitropoulos, IPE/DP, EPFL, PHB-Ecublens, 1015, Lausanne, Switzerland.

0932-0784 / 92 / 0100-0261 \$ 01.30/0. – Please order a reprint rather than making your own copy.



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland Lizenz.

Zum 01.01.2015 ist eine Anpassung der Lizenzbedingungen (Entfall der Creative Commons Lizenzbedingung „Keine Bearbeitung“) beabsichtigt, um eine Nachnutzung auch im Rahmen zukünftiger wissenschaftlicher Nutzungsformen zu ermöglichen.

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

On 01.01.2015 it is planned to change the License Conditions (the removal of the Creative Commons License condition “no derivative works”). This is to allow reuse in the area of future scientific usage.

structural phase transition is marked by a discontinuity in the slope of  $\nu_Q$  versus  $T$ . On deuteration,  $T_{c1}$  is slightly shifted to lower temperature, i.e. from 88 K to 86 K. The absolute values of  $\nu_Q$  show also only a small isotope effect, indicating that the influence of deuteration on the static average electric field gradient (efg) at the chlorine site is rather small. On the other hand, the

dynamical effects of deuteration are more pronounced as evidenced by the temperature dependence of the NQR relaxation rate near  $T_{c1}$  (see Figure 2).

Below  $T_{c1}$ , a shoulder appears, which is not directly related to the rotary soft mode as will be discussed further on. The ratio of the relaxation rate for the two chlorine isotopes is constant through the transition and very close to the ratio ( $\cong 1.6$ ) of the squares of their quadrupole moments  $Q$ , indicating that the relaxation mechanism is driven by small fluctuations of the quadrupole interaction.

### Discussion and Conclusions

The enhancement of the relaxation rate near  $T_{c1}$  is normally related to the anomalous temperature dependence of generalized unstable lattice modes near a structural ferro-distortive phase transition [3, 6]. Since the relaxation rates of the  $^{35,37}\text{Cl}$  isotopes scales as  $Q^2$  (see Fig. 2), a perturbative approach which relates the relaxation transition probabilities to the nuclear spin-phonon-scattering, is appropriate near  $T_{c1}$  (weak collision limit). This situation is to be compared to the one occurring at temperatures above 300 K, where an enhancement of  $T_1^{-1}$  is observed, due to sudden and slow  $90^\circ$  reorientations of the  $\text{TeCl}_6^{2-}$  octahedra. In this case the relaxation rates are identical for both chlorine isotopes, since the mechanism involves a change of the total quadrupole hamiltonian (strong

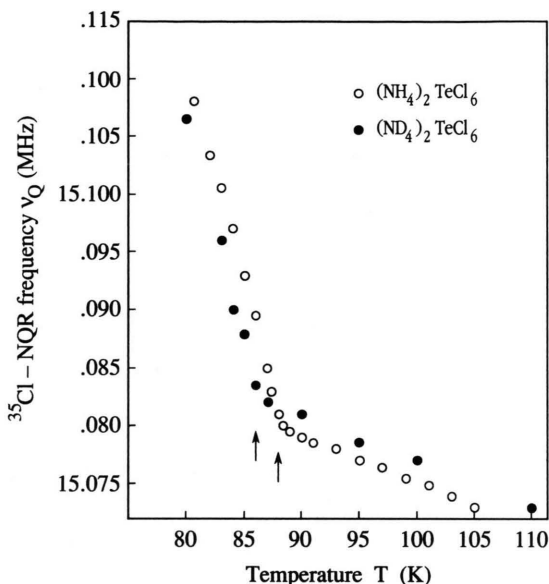


Fig. 1.  $^{35}\text{Cl}$ -NQR frequency  $\nu_Q$  vs. temperature in natural and deuterated hexatellurates. The arrows indicate the temperatures of the structural transition.

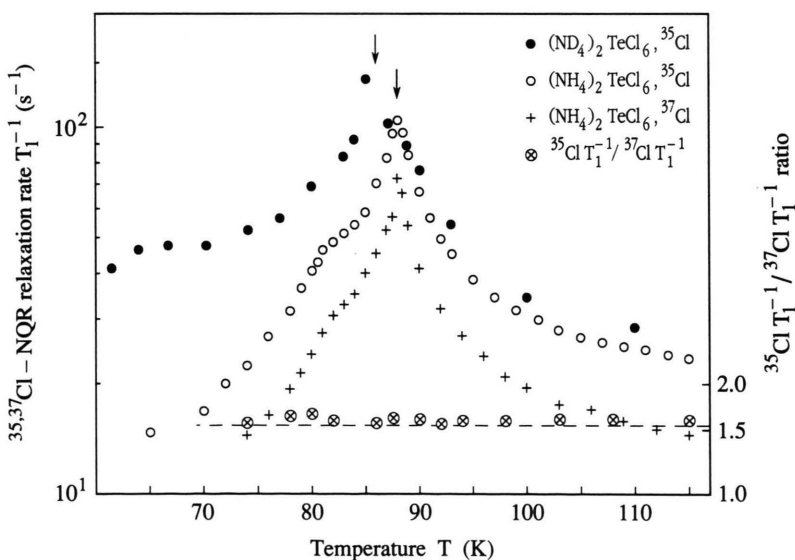


Fig. 2.  $^{35,37}\text{Cl}$ -NQR relaxation rate vs. temperature in hexatellurates. The arrows indicate the temperature of the structural transition. The dashed line corresponds to the ratio of the square of the quadrupole moments for the two Cl isotopes.

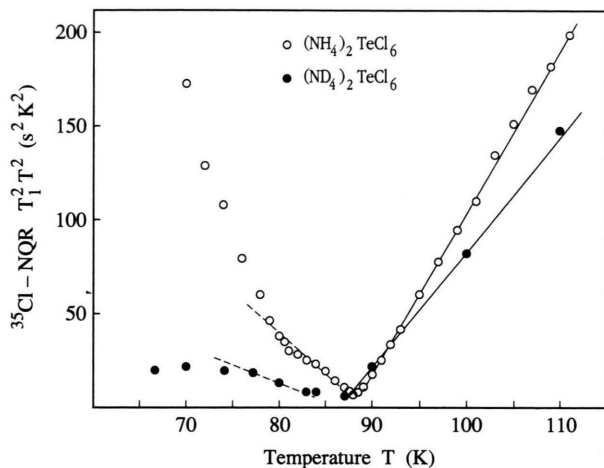


Fig. 3. Plot of  $(T_1 T)^2$  vs. temperature for the hexatellurates. The straight lines correspond to the fit to the critical behavior  $T_1^{-1} \propto (T - T_{c1})^{1/2}$ .

collision limit) leading to a relaxation rate  $T_1^{-1}$  which yields directly the reorientation rate  $\tau^{-1}$  [8].

The analysis of the experimental results in Fig. 3 indicates that the relaxation rate diverges at  $T \rightarrow T_{c1}$  according to a power law

$$T_1^{-1} = AT/(T - T_{c1})^{1/2}. \quad (1)$$

The coupling coefficient  $A$  is different in the cubic and the trigonal phases and depends on the degree of deuteration. The power law of (1) corresponds to the theoretical prediction obtained by considering a dominant relaxation mechanism due to a direct spin-phonon process in presence of overdamped generalized soft mode [10]. The square root singularity is characteristic of an isotropic interaction as one would expect in the present case, if the efg at the chlorine site were modulated by rotational vibrations of the  $\text{TeCl}_6^{2-}$  octahedra around each of the three equivalent pseudo-cubic  $[1, 0, 0]$  axes with equal probability. A generalized soft mode made up of linear combination of equal amplitude rotational vibrations around any of the  $[1, 0, 0]$  axes is consistent with the trigonal sym-

try of the frozen-in phase below  $T_{c1}$ . It should be emphasized that in the  $\text{K}_2\text{OsCl}_6$  salt, where the soft mode is only weakly damped, the divergence of  $T_1^{-1}$  at  $T_{c1}$  is found to obey a  $(T - T_{c1})^{-1}$  law as predicted for a Raman two-phonon relaxation mechanism appropriate in the case of underdamped soft modes [10, 11]. Thus one can conclude that the coupling between the  $\text{TeCl}_6^{2-}$  octahedra and the  $\text{NH}_4^+$  tetrahedra causes an enhanced damping of the rotary soft mode. This conclusion is strongly supported by recent Raman scattering measurements, showing an overdamping of the librational soft mode of the  $\text{TeCl}_6^{2-}$  octahedra near the critical region [12]. The values of the constant  $A$  in (1) as deduced from the reciprocal of the slopes of the straight lines in Fig. 3 are:  $A = 0.34 \text{ (sec}^{-1} \text{ K}^{-1/2})$  for  $(\text{NH}_4)_2\text{TeCl}_6$  and  $A = 0.4 \text{ (sec}^{-1} \text{ K}^{-1/2})$  for  $(\text{ND}_4)_2\text{TeCl}_6$ . If one assumes that the dominant effect on the coefficient  $A$  comes from the width  $\Gamma$  of the rotary soft mode,  $A \propto \Gamma$  [10], one can tentatively conclude that the effect of deuteration is on increasing the damping of the soft mode.

We consider now briefly the origin of the shoulder displayed in the  $T_1^{-1}$  versus  $T$  curve just below  $T_{c1}$  which is particularly pronounced for the deuterated compound. As it was argued previously [7], this shoulder is due to an extra contribution to the chlorine-NQR relaxation rate, associated with the modulation of the efg during the tumbling motion of the  $\text{NH}_4^+$  tetrahedra among three equivalent positions in a triangular shape potential. It is plausible, that the less pronounced shoulder observed in the natural salt is due to the same mechanism. The fact that this extra contribution to the relaxation is dependent upon deuteration both in magnitude and in the temperature position of the maximum, is in agreement with an isotope dependence of that triangle-shaped potential well.

#### Acknowledgements

The authors thank S. Müller for preparing the crystals used in this work.

- [1] R. L. Armstrong and H. M. van Driel, in *Adv. Nucl. Quadrupole Res.* 2, p. 179, Ed. J. A. S. Smith, Heyden 1975.
- [2] R. L. Armstrong, *Physics Reports* **57**, 343 (1980).
- [3] C. Dimitropoulos and J. Pelzl, *Z. Naturforsch.* **44a**, 109, (1989).
- [4] U. Kawald, S. Müller, J. Pelzl, and C. Dimitropoulos, *Sol. State Comm.* **67**, 239 (1988).
- [5] R. L. Armstrong, P. Dufort, and B. M. Powell, *Can. J. Phys.* **69**, 137 (1991).
- [6] S. Müller, U. Kawald, and J. Pelzl, in *Phonons 89*, Eds. S. Hunklinger, W. Ludwig, and G. Weiss, World Scientific, Vol. 2, p. 1120, Singapore 1990.
- [7] C. Dimitropoulos, J. Pelzl, and F. Borsa, *Phys. Rev. B* **41**, 3914 (1990).
- [8] Y. Furukawa *et al.*, *Bull. Chem. Soc. Japan* **50**, 1927 (1977).
- [9] J. J. van der Klink and C. Dimitropoulos, *J. Molec. Structure* **58**, 359 (1980) and references therein.
- [10] G. Bonera, F. Borsa, and A. Rigamonti, *Phys. Rev. B* **2**, 2784 (1970).
- [11] R. L. Armstrong, M. E. Ramia, and R. M. Morra, *J. Phys. C* **19**, 4363 (1986).
- [12] S. Müller, Ph. D. thesis, Ruhr University Bochum, FRG (1991).